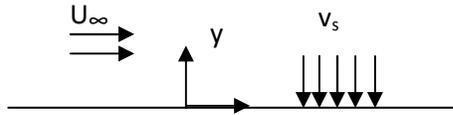


Problem 2.4: Solution



Assume steady, incompressible, constant-property, laminar flow. Also, assume negligible gravity effect. For an infinitely large plane no dependence on x is possible, therefore

$$\rho \left(\frac{du}{dt} + u \frac{du}{dx} + v \frac{du}{dy} \right) = - \frac{dP}{dx} + \mu \left(\frac{d^2u}{dx^2} + \frac{d^2u}{dy^2} \right)$$

$$\begin{matrix} \swarrow & \swarrow & \swarrow & \swarrow \\ 0 & 0 & 0 & 0 \end{matrix}$$

$$v \frac{du}{dy} - \nu \frac{d^2u}{dy^2} = 0 \tag{1}$$

The y -momentum gives

$$\rho \left(\frac{dv}{dt} + u \frac{dv}{dx} + v \frac{dv}{dy} \right) = - \frac{dP}{dy} + \mu \left(\frac{d^2v}{dx^2} + \frac{d^2v}{dy^2} \right)$$

$$\begin{matrix} \swarrow & \swarrow & \swarrow & \swarrow \\ 0 & 0 & 0 & 0 \end{matrix}$$

$$v \frac{dv}{dy} - \nu \frac{d^2v}{dy^2} = 0 \tag{2}$$

However, mass conservation says,

$$\frac{du}{dx} + \frac{dv}{dy} = 0 \rightarrow \frac{dv}{dy} = 0$$

$$0$$

$$\rightarrow v = \text{constant} = -v_s \tag{3}$$

Evidently, $v = v_s$ everywhere, and that satisfies Eq. (2).

Eqn. (1) can now be written as

$$\frac{d^2u}{dy^2} + \frac{v_s}{\nu} \frac{du}{dy} = 0 \tag{4}$$

$$\text{At } y = 0, u = 0$$

$$\text{At } y \rightarrow \infty, u = U_\infty$$

The solution will be

$$u = U_\infty \left[1 - \exp \left(-\frac{v_s}{\nu} y \right) \right] \tag{5}$$

(Note that, by our definition, $v_s > 0$ for suction at the wall.)

If we define the edge of the boundary layer as the location where $\frac{u}{U_\infty} = \zeta^*$, then

$$\zeta^* = \left[1 - \exp\left(-\frac{v_s}{\nu} \delta\right) \right] \quad (6)$$

For $v_s = 0.1 \text{ m/s}$, and assuming $\nu = 8.6 \times 10^{-7} \text{ m}^2/\text{s}$ for water, and $\zeta^* = 0.999$, then

$$\delta \approx 6 \times 10^{-5} \text{ m} \quad (7)$$

For blowing, $v_s < 0$, and since $\zeta^* < 1$, will set,

$$\delta = -\frac{\nu}{v_s} \ln(1 - \zeta^*) < 0 \quad (8)$$

(No acceptable solution)

a) In steady-state, neglecting the viscous dissipation, the energy equation becomes

$$u \frac{dT}{dx} - v_s \frac{dT}{dy} = \alpha \frac{d^2 T}{dy^2}$$

$$\begin{aligned} & 0 \\ \rightarrow & \frac{dT}{dy} + \frac{\alpha}{v_s} \frac{d^2 T}{dy^2} = 0 \end{aligned} \quad (9)$$

$$\text{At } y = 0, T = T_s \quad (10)$$

$$\text{At } y \rightarrow \infty, T \rightarrow \infty \quad (11)$$

The solution will be

$$\frac{T - T_\infty}{T_s - T_\infty} = 1 - \exp\left(-\frac{v_s}{\alpha} y\right) \quad (12)$$

For air, with $\nu = .158 \times 10^{-4} \text{ m}^2/\text{s}$, will have for the suction case,

$$\delta = 0.00109 \text{ m}$$

$$\delta_3 = \int_0^\infty \frac{\rho u}{\rho_\infty U_\infty} \left(1 - \frac{u^2}{U_\infty^2}\right) dy$$

$$= \left(\frac{U_\infty}{\nu x}\right)^{-1/2} \int_0^5 \left(\sin \frac{\pi}{10} \eta - \sin^3 \frac{\pi}{10} \eta\right) d\eta$$

$$= \left(\frac{U_\infty}{\nu x}\right)^{-1/2} \left(\frac{10}{3\pi}\right)$$